

# **Axial U(1) symmetry at high temperature in 2-flavor lattice QCD**

**Presenter:**

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**from JLQCD Collaboration:**

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# Outline of this talk

1. Motivation
2. Axial  $U(1)$  symmetry and Dirac spectrum
3. Results
  - 3-1:  $U(1)_A$  susceptibility at finite  $T$
  - (3-2: Topological susceptibility)  $\Rightarrow$  Next talk by [Y. Aoki](#)
4. Summary

Related talks from JLQCD:

[H. Fukaya](#) (Thu. plenary) – Summary of finite  $T$

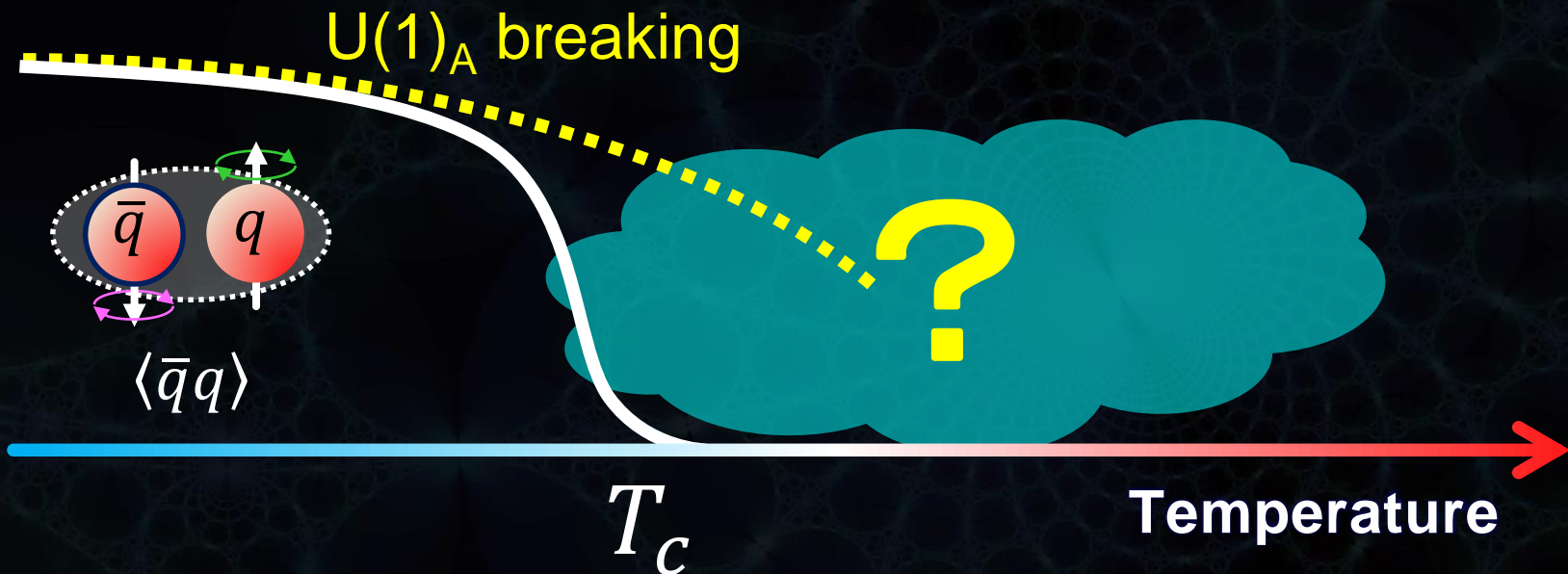
[H. Fukaya](#) (Thu.) – Topological susceptibility at zero  $T$

[C. Rohrhofer](#) (Wed.) – Mesonic correlator at finite  $T$

# Motivation: Is $U(1)_A$ symmetry restored above $T_c$ ?

- For  $T > T_c$ , chiral symmetry breaking by  $\langle \bar{q}q \rangle$  is restored  
 $\Rightarrow$  How about  $U(1)_A$  symmetry?

$$\Delta_{\pi-\delta} = \int_0^\infty d^4x [\pi^a(x)\pi^a(x) - \delta^a(x)\delta^a(x)]$$



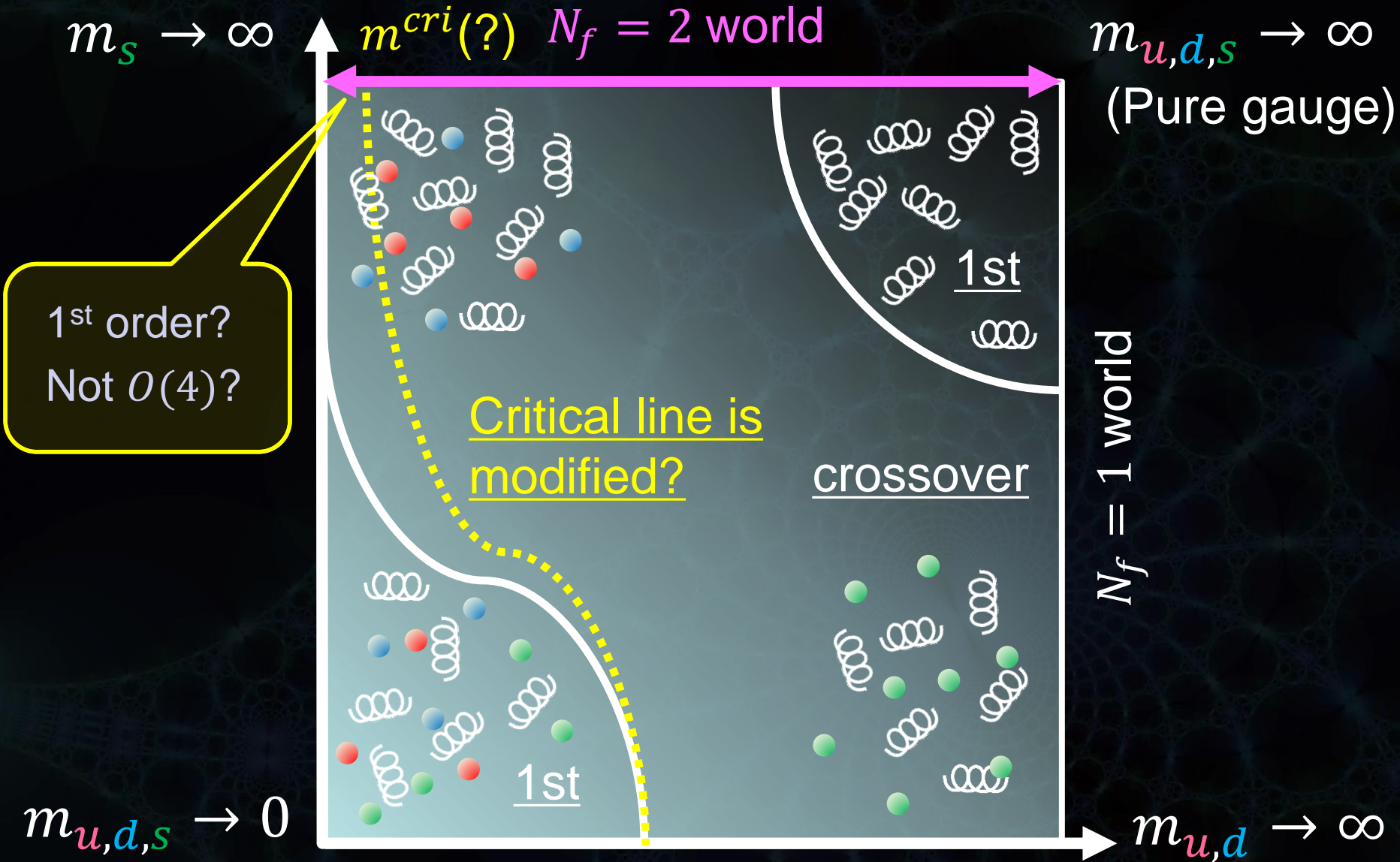
# Motivation: Is $U(1)_A$ symmetry restored above $T_c$ ?

- For  $T \gg T_c$ ,  $U(1)_A$  symmetry will be restored because of dilute instanton gas picture (Gross-Pisarski-Yaffe, 1981)
- HotQCD (DW, 2012) **broken**
- JLQCD (topology fixed overlap, 2013) **restored**
- TWQCD (optimal DW, 2013) **restored**
- LLNL/RBC (DW, 2014) **broken**
- Dick et al. (overlap on HISQ, 2015) **broken**
- JLQCD (reweighted overlap from DW, 2016) **restored**

⇒ It is still an unsettled problem

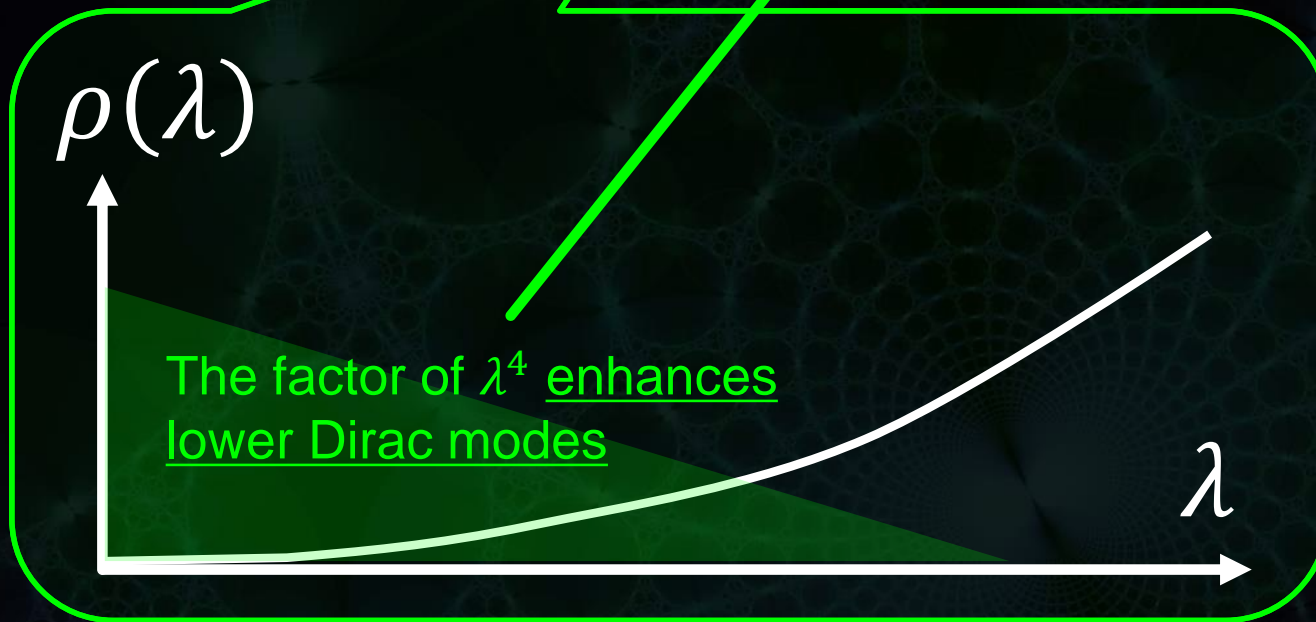
If  $U(1)_A$  is restored...

# the Columbia plot is modified?



# $U(1)_A$ susceptibility from Dirac spectra

$$\Delta_{\pi-\delta} = \int_0^\infty d\lambda \rho(\lambda) \frac{2m^2}{(\lambda^2 + m^2)^2}$$

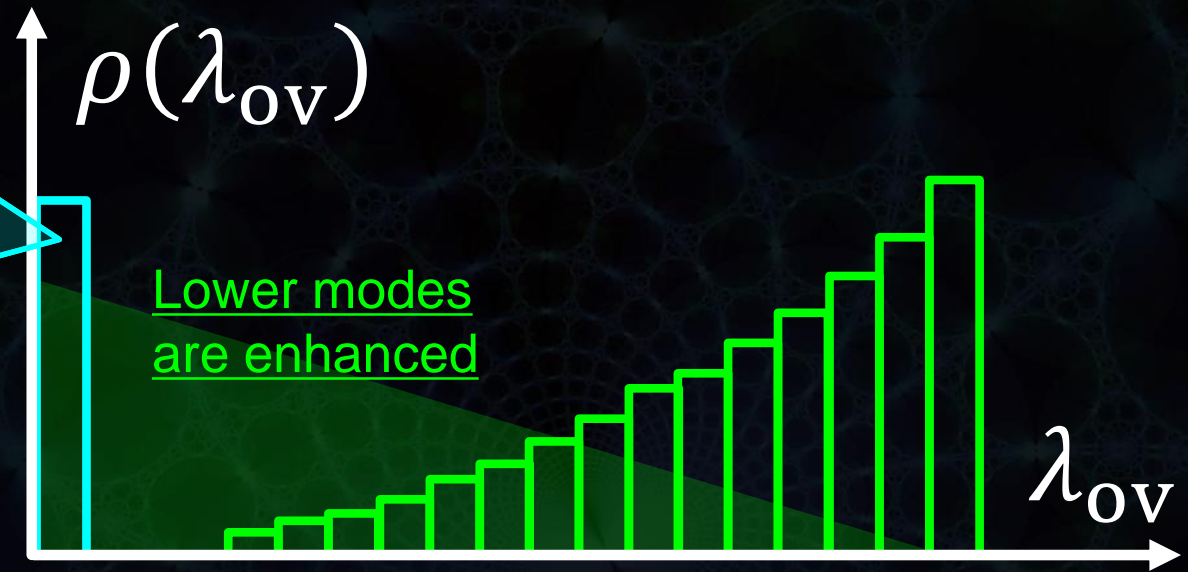


Cf.) Banks-Casher relation:  $\langle \bar{q}q \rangle = \lim_{m \rightarrow 0} \int_0^\infty d\lambda \rho(\lambda) \frac{2m}{\lambda^2 + m^2}$

# Observables: Dirac spectra on the lattice

$$\Delta_{\pi-\delta} = \int_0^\infty d\lambda \rho(\lambda) \frac{2m^2}{(\lambda^2 + m^2)^2} \quad \Rightarrow \quad \Delta_{\pi-\delta}^{\text{ov}} \equiv \frac{1}{V(1-m^2)^2} \sum_i \frac{2m^2(1-\lambda_{\text{ov}}^{(i)4})^2}{\lambda_{\text{ov}}^{(i)4}}$$

Practically, (physical) zero-modes will be dominant  
 $\Rightarrow$  instanton contributions in finite box



Practical observable:  $\bar{\Delta}_{\pi-\delta}^{\text{ov}} \equiv \frac{1}{V(1-m^2)^2} \sum_{i>0\text{-mode}} \frac{2m^2(1-\lambda_{\text{ov}}^{(i)4})^2}{\lambda_{\text{ov}}^{(i)4}}$   
 Subtracted zero-modes

(Instanton contribution becomes zero at  $V \rightarrow \infty$ :  $\bar{\Delta}_{0\text{-mode}}^{\text{ov}} = \frac{2N_0}{Vm^2} (\propto 1/\sqrt{V})$ )

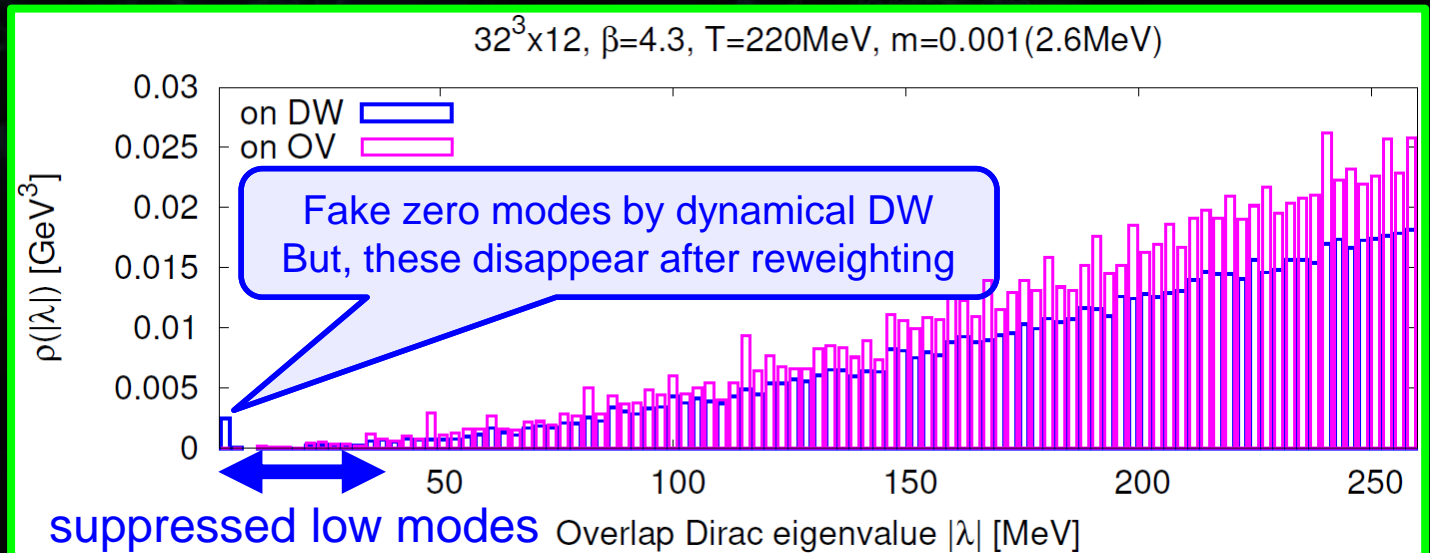
# Project of $U(1)_A$ at finite $T$ by JLQCD collaboration

	Valence/Sea quarks	Setup
G. Cossu et al. PRD87 (2013)	OV on OV (Topology fixed sector)	
A. Tomiya et al. [arXiv:1612.01908]	DW on DW OV on DW <u>OV on (reweighted) OV</u>	$1/a=1.7\text{GeV}$
<u>This work</u>	OV on DW <u>OV on (reweighted) OV</u>	$1/a=2.6\text{GeV}$ (More fine)

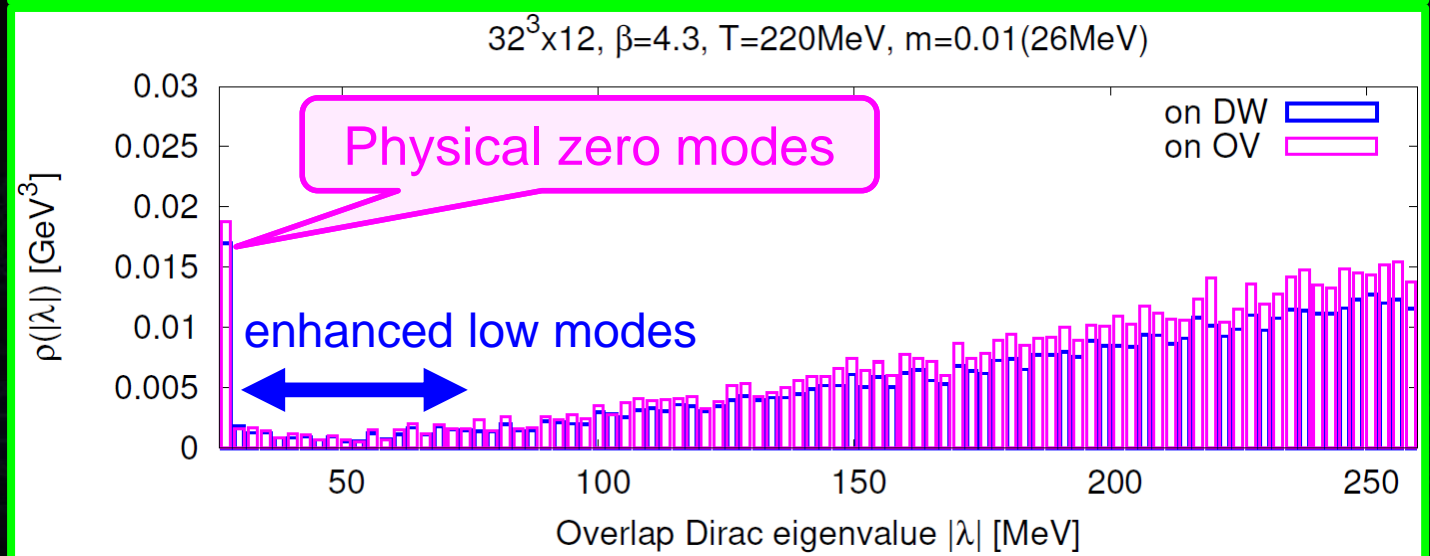


# Results (overlap Dirac spectra)

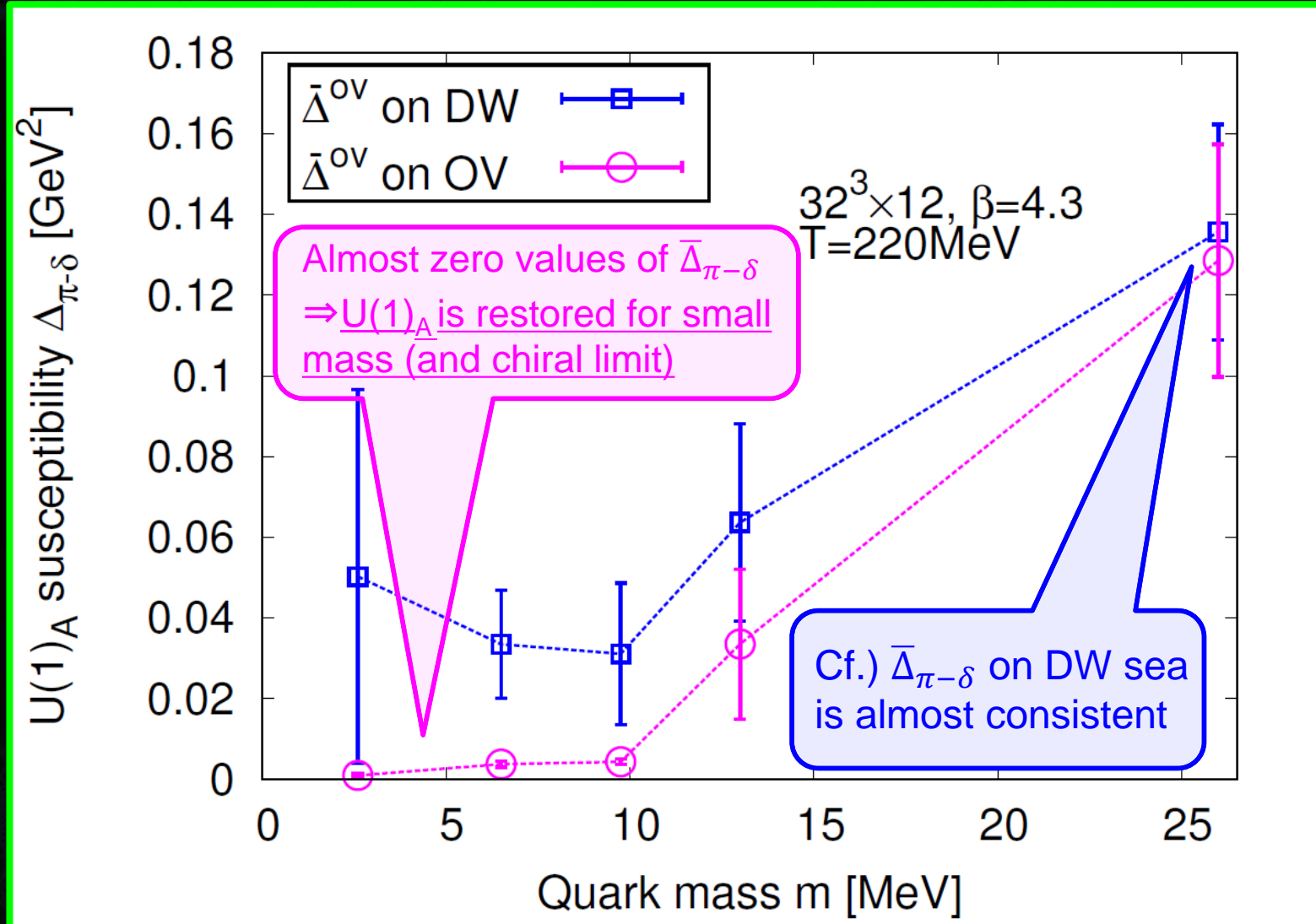
$m=2.6\text{MeV}$



$m=26\text{MeV}$

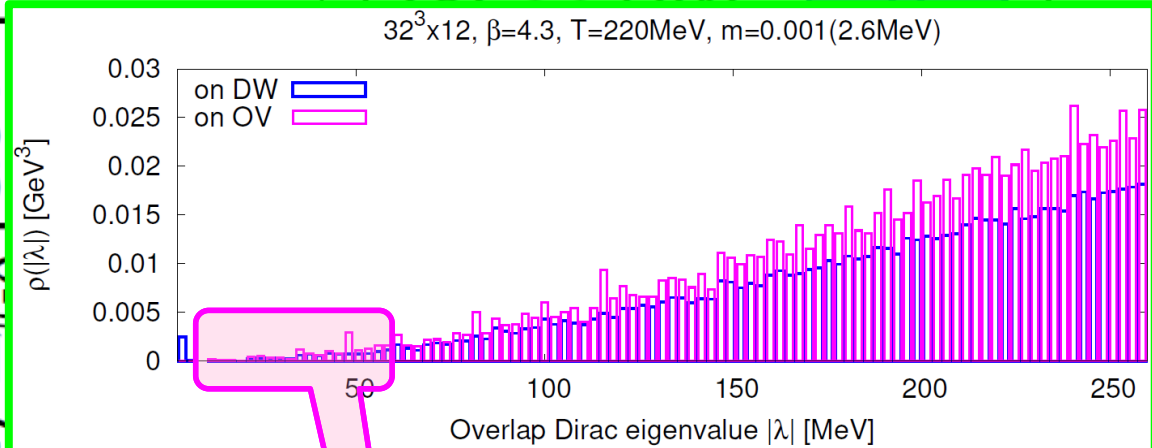


# $U(1)_A$ symmetry at $T = 220\text{MeV}$

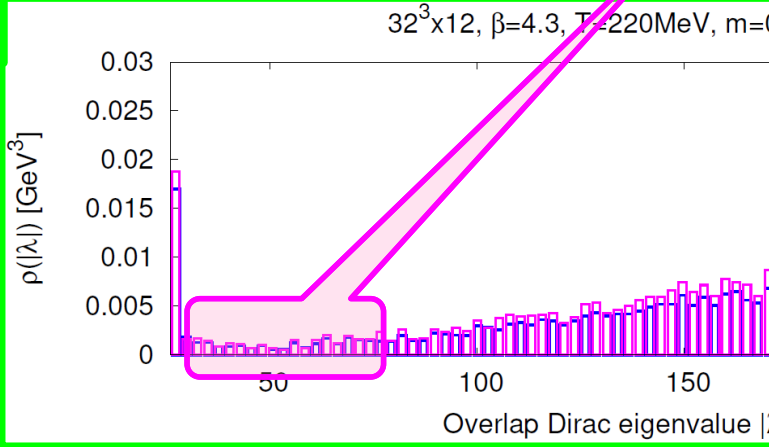


0.18

$= 0V$



32<sup>3</sup> × 12, β=4.3  
T=220MeV



⇒ Finite  $\bar{\Delta}_{\pi-\delta}$  comes from enhancement of lower Dirac modes

0.1

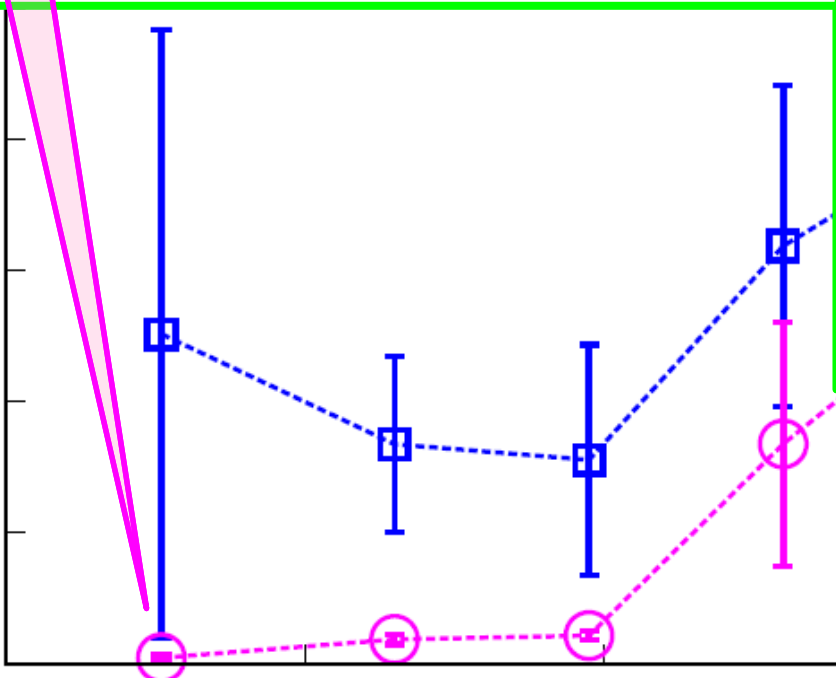
0.08

0.06

0.04

0.02

0



⇒ Zero  $\bar{\Delta}_{\pi-\delta}$  comes from suppression of lower Dirac modes

10

15

20

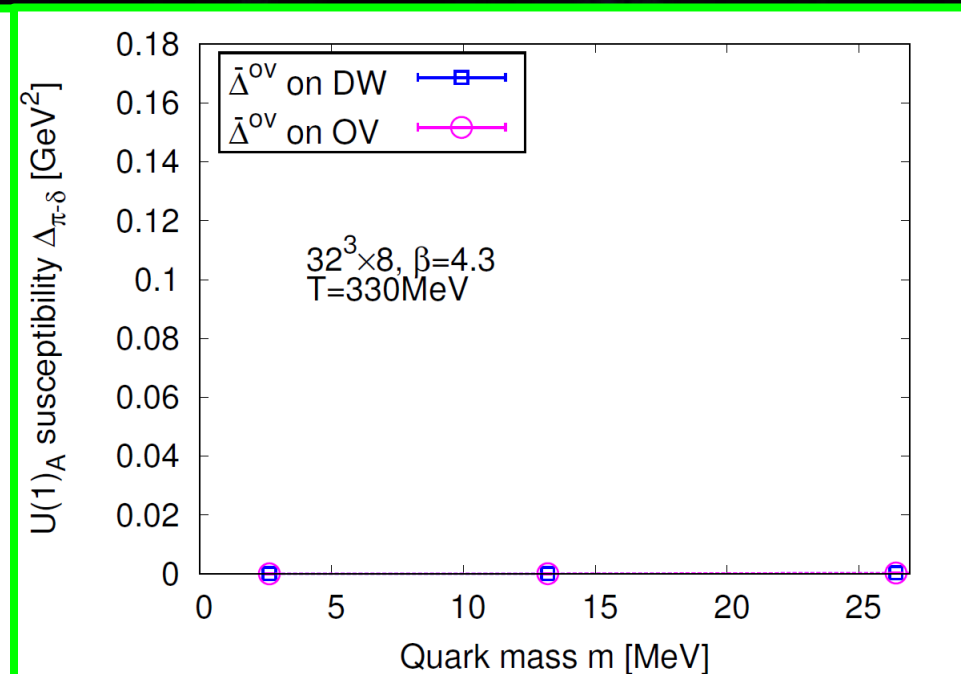
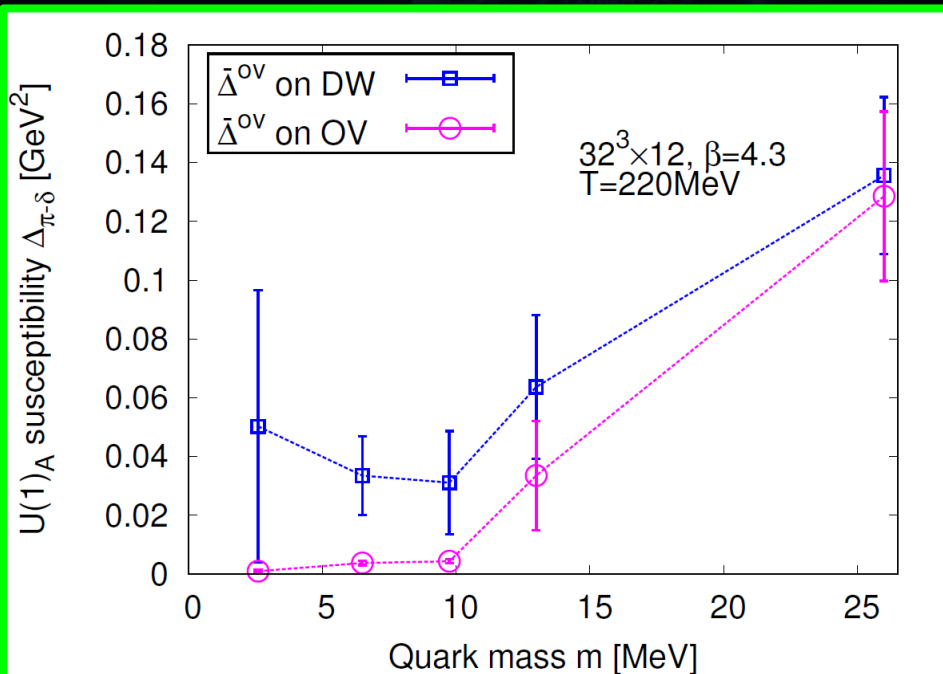
25

Quark mass m [MeV]

# Temperature dependence (220 or 330MeV)

T=220MeV

T=330MeV

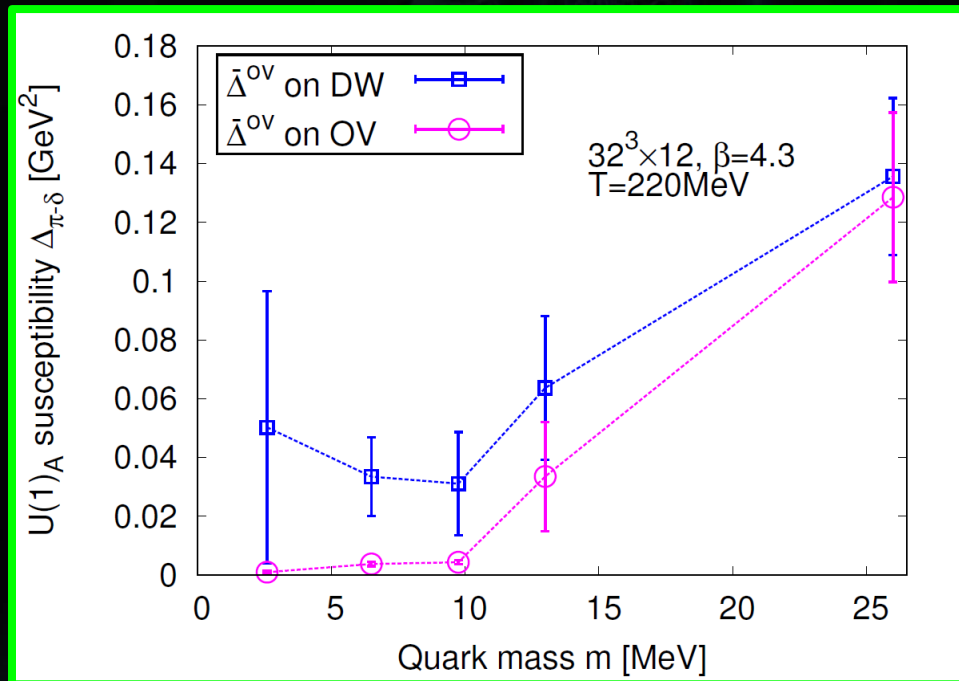


⇒ With increasing  $T$ ,  $U(1)_A$  symmetric region becomes wider

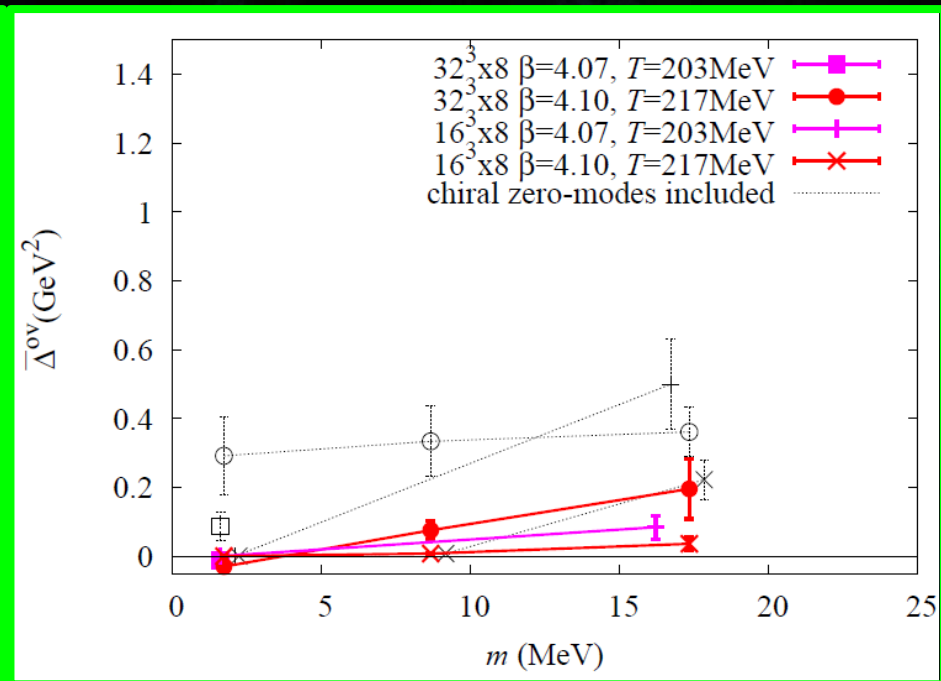
# Comparison with Coarse lattice

JLQCD, preliminary (2017)

A. Tomiya et al. (JLQCD) [arXiv:1612.01908]



$1/a=2.6\text{GeV}$



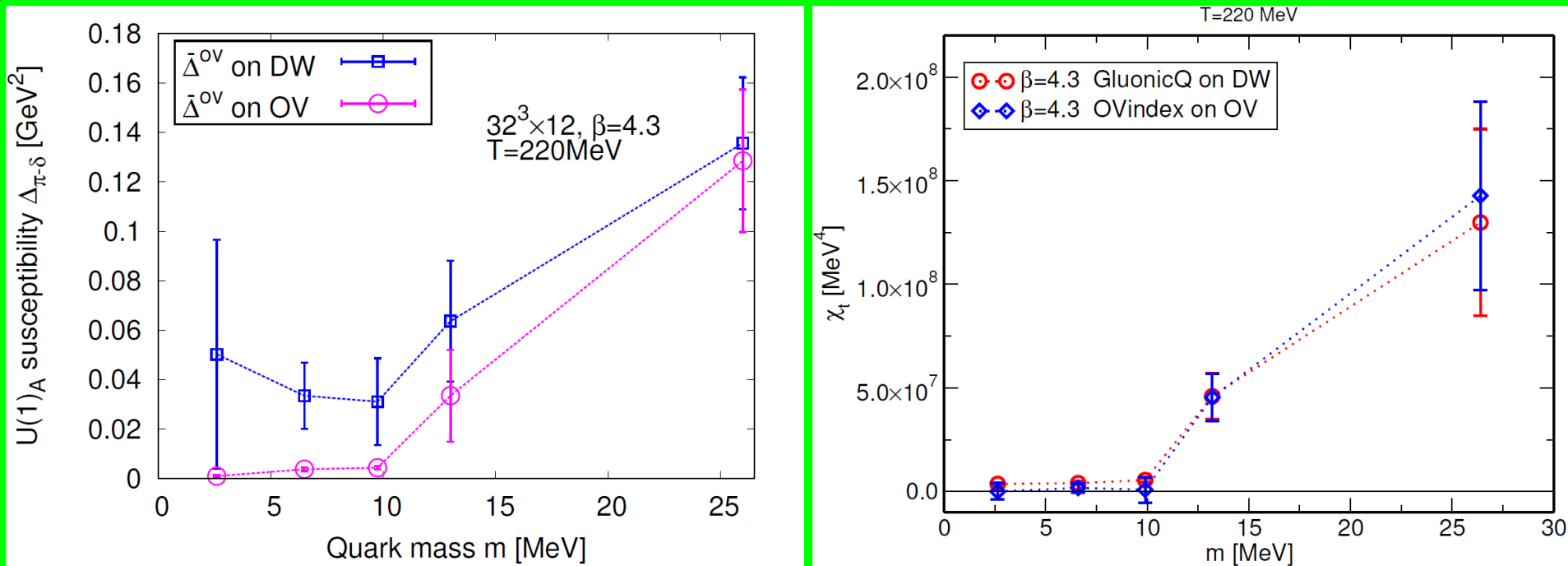
$1/a=1.7\text{GeV}$

⇒ Results from two lattice are consistent

# Comparison with Topological susceptibility

- $Q_t$  is related to Dirac operator by the index theorem:

$$Q_t = n_+ - n_- = \text{index}(D)$$

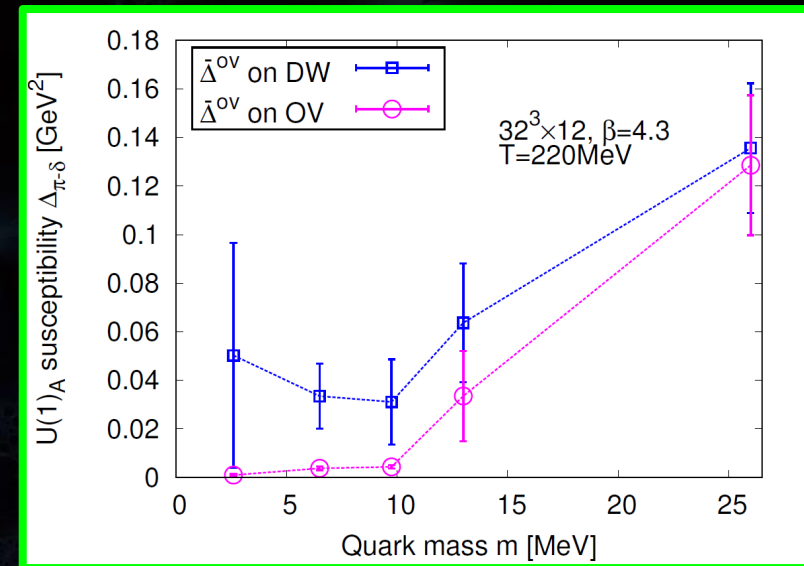


The behavior seems to be similar? [⇒ Next talk by Y. Aoki](#)

# Summary and Outlook

- $U(1)_A$  susceptibility at high  $T$  for  $N_f = 2$  was simulated from OV Dirac spectra on the chiral sea quarks

⇒ It seems to be restored in the chiral limit for  $T=220-330\text{MeV}$



- Relation to topological susceptibility ⇒ Next talk
- Functional form of  $\Delta(m)$  (powers of  $m$  ...?)
- Going to near  $T_c$  (critical exponent for chiral transition?)
- Going to  $N_f = 2 + 1$  sector

# Backup