

# Spectroscopy of the BSM sextet model

Outline

Challenges

Mixed Action

Preliminary results

Conclusion

Chik Him (Ricky) Wong

Lattice Higgs Collaboration (L<sub>at</sub>HC):

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LATTICE 2017

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- Challenges of lattice analysis of nearly conformal BSM model
- Mixed action analysis with application of gradient flow
- Preliminary results
- Conclusion

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- Case study of nearly conformal model: Sextet model  $SU(3)$  with  $N_f = 2$  in two-index symmetric representation
- Focus of this talk: Hadron Spectroscopy study
- No IRFP is found despite small  $\beta$  function [more in Kieran Holland's talk]  
⇒ still chirally broken
- Is  $\chi$ PT applicable? Dynamics are very different from QCD
- A low-lying  $0^{++}$  scalar exists  
⇒ Linear- $\sigma$  model or Dilaton-inspired extended theory? (Julius Kuti's talk)
- Expensive to get lower in  $m$  where  $0^{++}$  scalar may decouple

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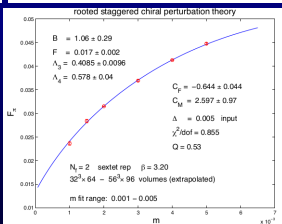
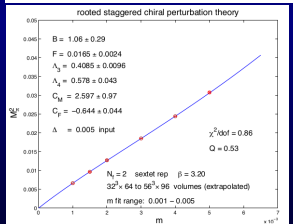
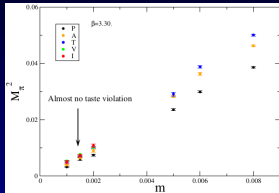
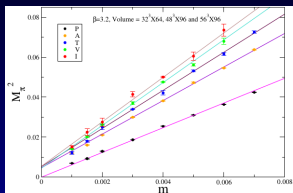


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- Staggered formalism:  
Taste breaking structure is different from QCD



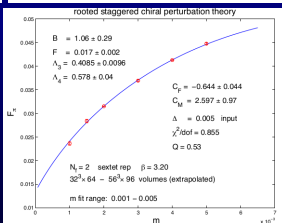
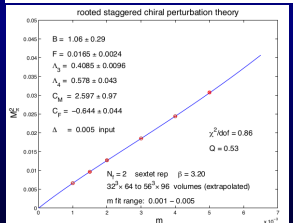
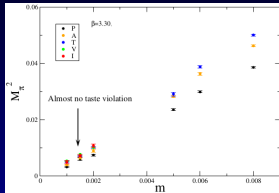
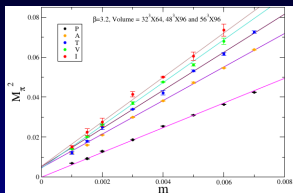
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- Goal: Ease the analysis by isolating taste breaking effects  
 $\Rightarrow$  Mixed action analysis

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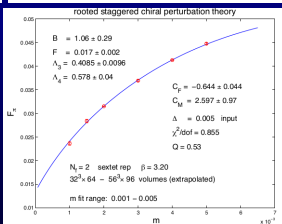
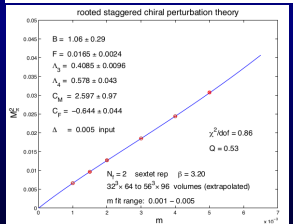
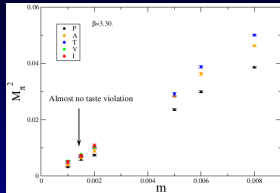
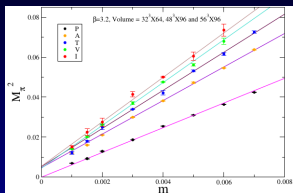
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- **Mixed action setup: Valence fermion action is different from sea fermion action**
- $\chi$ PT  $\Rightarrow$  MA $\chi$ PT
- If the valence fermion action (almost) respects taste symmetry,
  - The valence Goldstone spectrum  $M_{\pi, V}$  has no detectable taste splitting
  - Taste breaking effect from the sea fermion action is messaged in the mixed meson  $M_{\pi, S}$
- Flowed action as the valence action :
  - Apply a large number of small stout smearing steps to the original links, equivalent to a Wilson flow of  $t_0 a^2$
  - Staggered action defined by the “flowed” links with a valence quark mass  $m_V$

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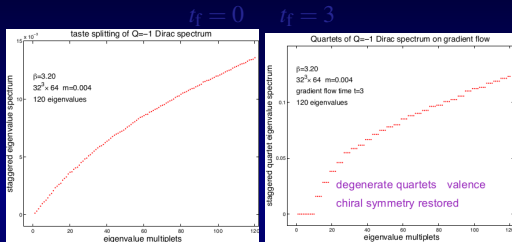
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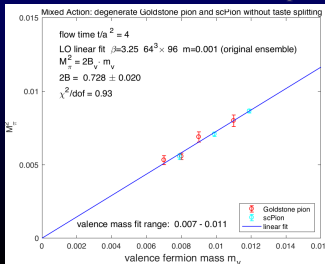
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## • Flowed action as the valence action :

- Wilson flow reduces cutoff effects, approximately restores the taste symmetry



- Valence Goldstone and valence non-Goldstone pions degenerate



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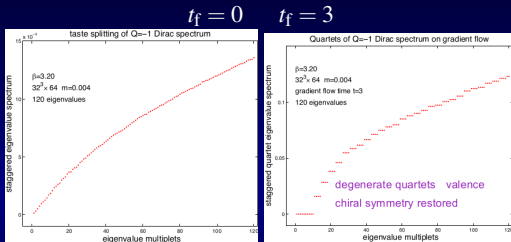
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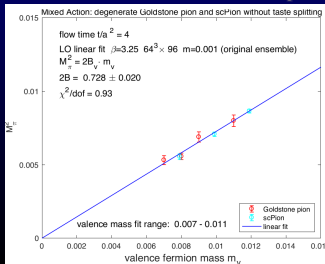
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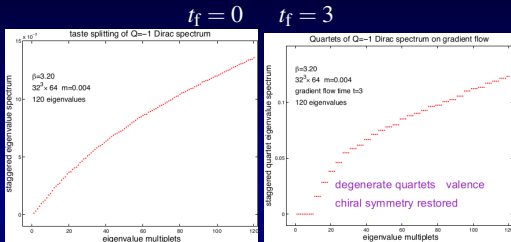
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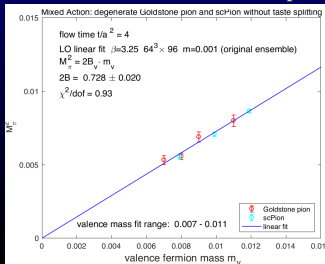
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- **Additional advantage:**
  - Wilson flow smoothens the links  $\Rightarrow$  Faster inversions
  - Hadron correlators, especially  $0^{++}$  scalar, becomes more affordable
  - Going lower in  $m_V$  is easier

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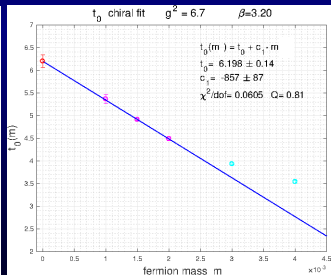
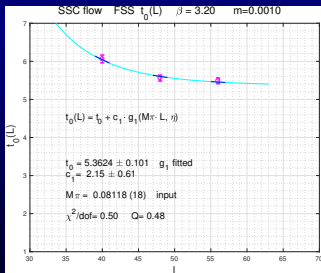
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- Measure quantities  $X$ , e.g.  $F_\pi$ ,  $M_X$  on gradient flowed links.  
 $\sqrt{8t_f} M_\pi \ll 1$  is required, and  $p$ -regime is assumed.
- Define a scale  $t_0$  :  $g^2|_{t_0} = c$  (6.7 here)
- Determine the chiral limit of the infinite volume extrapolated  $t_0$  at each  $\beta$  [ Kieran Holland's talk]



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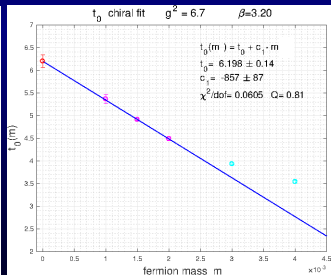
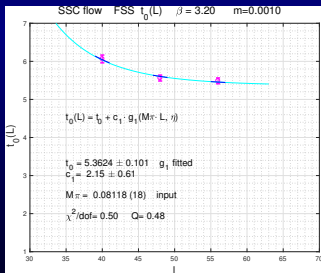
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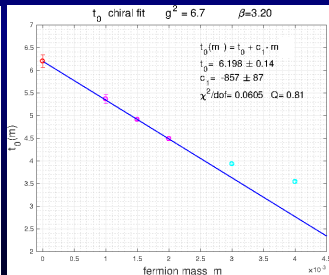
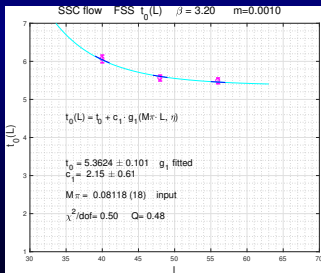
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- Assuming MA $\chi$ PT, the relation between  $X_{\text{NLO}}$  and  $X_{\text{LO}}$  in the infinite-volume limit can be worked out. E.g.

$$\frac{M_{\pi,vv,\text{NLO}}^2}{M_{\pi,vv,\text{LO}}^2} = \left( 1 + \frac{1}{32\pi^2 F_\pi^2} \left[ (M_{\pi,vv,\text{LO}}^2 - M_{\pi,ss,\text{LO}}^2) \right. \right. \\ \left. \left. + (2M_{\pi,vv,\text{LO}}^2 - M_{\pi,ss,\text{LO}}^2) \ln \frac{M_{\pi,vv,\text{LO}}^2}{\mu^2} \right] \right. \\ \left. - \frac{8}{F_\pi} \left[ (L_5 - 2L_8)M_{\pi,vv,\text{LO}}^2 + 2(L_4 - 2L_6)M_{\pi,ss,\text{LO}}^2 \right] + C a^2 \right), \\ \frac{F_{\pi,vv,\text{NLO}}^2}{F_\pi} = \left( 1 - \frac{M_{\pi,vs,\text{LO}}^2}{16\pi^2 F_\pi^2} \ln \frac{M_{\pi,vs,\text{LO}}^2}{\mu^2} + \frac{4}{F_\pi^2} \left[ L_5 M_{\pi,vv,\text{LO}}^2 + 2L_4 M_{\pi,ss,\text{LO}}^2 \right] \right. \\ \left. + D a^2 \right),$$

where  $M_{\pi,vv,\text{LO}}^2 = 2Bm_v$ ,  $M_{\pi,ss,\text{LO}}^2 = 2Bm_s$ ,  $M_{\pi,vs,\text{LO}}^2 = B(m_v + m_s)$

- Scaling with  $t_0$  to obtain dimensionless quantities  $\bar{y} \equiv \sqrt{t_0} y$ , these relations can be parameterized as

$$\bar{M}_{\pi, \nu\nu, \text{NLO}}^2 = \bar{m}_\nu \left( 2\bar{B} + b_1 \bar{m}_\nu + b_2 \bar{m}_s + b_3 (2\bar{m}_\nu - \bar{m}_s) \ln \bar{m}_\nu + b_4 t_0^{-1} a^2 \right)$$

$$\bar{F}_{\pi, \nu\nu, \text{NLO}} = \bar{F}_\pi + c_1 \bar{m}_\nu + c_2 \bar{m}_s + c_3 (\bar{m}_\nu + \bar{m}_s) \ln \bar{m}_\nu + c_4 t_0^{-1} a^2$$

- In Mixed action, chiral limit and continuum limit cannot be separated  
 $\Rightarrow$  Simultaneous fit of all available  $L^3 \times T$ 's,  $\beta$ 's (with corresponding  $t_0$ 's),  $m_s$  and  $m_\nu$ 's to obtain  $\bar{X}$ 's in the infinite-volume, chiral and continuum limit
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$$\bar{M}_{\pi, \nu\nu, \text{NLO}}^2 = \bar{m}_\nu \left( 2\bar{B} + b_1 \bar{m}_\nu + b_2 \bar{m}_s + b_3 (2\bar{m}_\nu - \bar{m}_s) \ln \bar{m}_\nu + b_4 t_0^{-1} a^2 \right)$$

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# Preliminary results

Spectroscopy of  
the BSM sextet  
model

Chik Him (Ricky)  
Wong

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- **Sea Action:**  
Tree-level Symanzik-Improved gauge action with 2-step  $\rho = 0.15$  stout-smearred Staggered  $N_f = 2$  SU(3) Sextet fermion, sea mass  $m_s$
- **Valence Action:**  
Staggered fermion action with  $t_f = 4$  (80 steps of  $\rho = 0.05$  stout smeared) gauge, valence mass  $m_v$
- **Available data:**

$\beta$	$m_s$	$m_v$	$L^3 \times T$
3.20	0.001	0.006,0.008,0.010	$40^3 \times 80, 48^3 \times 96, 56^3 \times 96$
	0.0015	0.006,0.008,0.010	$56^3 \times 96$
	0.002	0.006,0.008,0.010	$56^3 \times 96$
3.25	0.001	0.006,0.008,0.010	$48^3 \times 96, 56^3 \times 96, 64^3 \times 96$
3.30	0.001	0.006,0.008,0.010	$48^3 \times 96, 56^3 \times 96, 64^3 \times 96$

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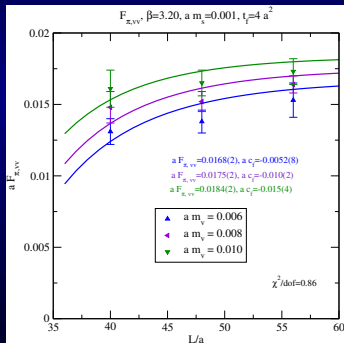
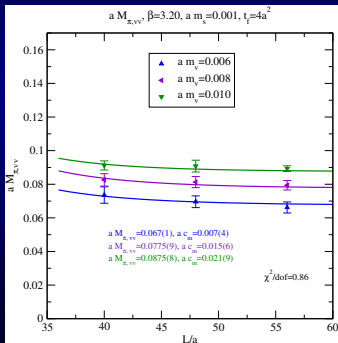
Mixed Action

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- Preliminary results of global fit: Finite Size Scaling part

$$M_{\pi, \nu\nu}(L) = M_{\pi, \nu\nu} + c_m g_1(M_{\pi, \nu\nu} L), \quad F_{\pi, \nu\nu}(L) = F_{\pi, \nu\nu} + c_f g_1(M_{\pi, \nu\nu} L)$$



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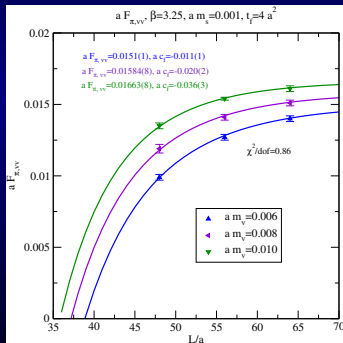
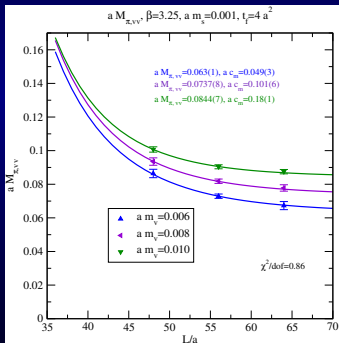
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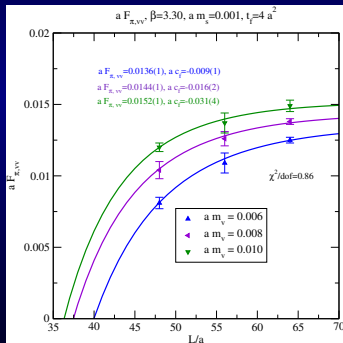
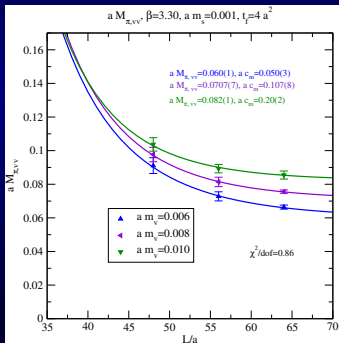
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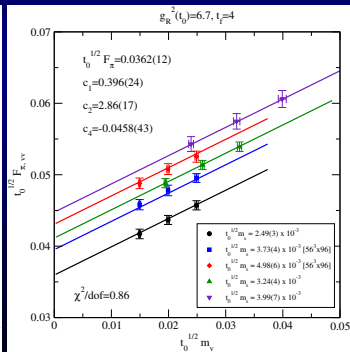
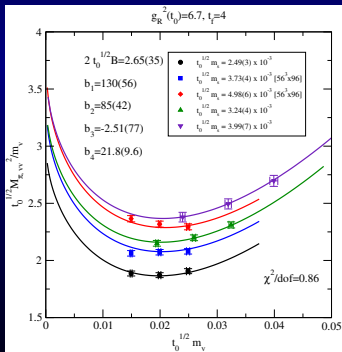
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$$\bar{M}_{\pi, vv, \text{NLO}}^2 = \bar{m}_v (2\bar{B} + b_1\bar{m}_v + b_2\bar{m}_s + b_3(2\bar{m}_v - \bar{m}_s) \ln \bar{m}_v + b_4 t_0^{-1} a^2)$$

$$\bar{F}_{\pi, vv, \text{NLO}} = \bar{F}_{\pi} + c_1\bar{m}_v + c_2\bar{m}_s + c_3(\bar{m}_v + \bar{m}_s) \ln \bar{m}_v + c_4 t_0^{-1} a^2$$

( $c_3$  is not detectable, hence omitted in the fit)





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- **Hadron Spectroscopy on nearly conformal models is challenging**
  - Low-lying  $0^{++}$  scalar interacts with pion dynamics  $\Rightarrow$   $\chi$ PT needs modification
  - Taste breaking structure for Staggered simulations is different from QCD
- **Mixed action analysis is explored**
  - Wilson-flowed links used as the valence action
  - Flowed action restores taste symmetry and allows cheaper inversions
  - Preliminary results is promising, demonstrating the viability of the approach
- **More comprehensive analysis is needed. E.g. Errors of  $t_0$ , Non-unitarity effects, the effects of low-lying  $0^{++}$  scalar**

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