

# Domain Wall Fermion QCD with the Exact One Flavor Algorithm

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The 35th International Symposium on Lattice Field Theory (Granada, Spain)

June 20th, 2017

Talk based on: C. Jung, C. Kelly, R.D. Mawhinney, and D.J. Murphy, *Domain Wall Fermion QCD with the Exact One Flavor Algorithm* [arXiv:1706.05843]

## Motivation

- Standard algorithm for lattice ensemble generation is hybrid Monte Carlo (HMC)
- Dynamical quark effects are described by determinant of fermion matrix

$$\langle \mathcal{O} \rangle = \frac{1}{\mathcal{Z}} \int \mathcal{D}U \mathcal{O}[U] \det(\mathcal{M}[U]) e^{-S[U]}$$

- Typically choose Hermitian  $\mathcal{M} = D^\dagger D$ , where  $D$  is the lattice Dirac operator, but  $\det(D^\dagger D)$  describes two degenerate quark flavors with same mass
- Take square root for one flavor simulation  $\implies$  rational HMC (RHMC)
- Multishift CG makes RHMC practical, but...
  - ▶ Generally more expensive than ordinary CG, due to extra linear algebra at each iteration
  - ▶ Can't use nonzero initial guesses, so techniques such as restarted solvers and forecasted initial guesses are not applicable
  - ▶ Empirically find Hasenbusch mass preconditioning to be prohibitively expensive
- For RBC/UKQCD  $\Delta I = 1/2 K \rightarrow \pi\pi$  calculation with  $G$ -parity boundary conditions  $\det(D^\dagger D)$  describes four flavors, and RHMC is needed for light quarks as well
- **Goal:** Use EOFA to accelerate  $K \rightarrow \pi\pi$  ensemble generation
  - ▶ New preconditioning scheme for EOFA Dirac equation
  - ▶ Timing benchmarks on production ensembles
- More detail in our paper!

## The Exact One Flavor Algorithm (EOFA)

- RHMC: Form rational approximation  $x^{1/2} \simeq \alpha_0 + \sum_k \alpha_k / (x + \beta_k)$  and compute

$$\det \left[ \frac{D(m_1)}{D(m_2)} \right] = \left\{ \det \left[ \frac{D^\dagger D(m_1)}{D^\dagger D(m_2)} \right] \right\}^{1/2}$$

Multishift CG allows  $D^\dagger D + \beta_k$  to be inverted for all  $k$  simultaneously

- EOFA: Use Schur determinant identity applied to spin structure of  $D$  to factorize

$$\det \left[ \frac{D(m_1)}{D(m_2)} \right] = \frac{1}{\det(H_1)} \cdot \frac{1}{\det(H_2)}$$

with  $H_1$  and  $H_2$  Hermitian and positive-definite [TWQCD, arXiv:1403.1683]

- No rational approximation, so no need for multishift CG
- Expect [TWQCD, arXiv:1412.0819]:
  - ▶ Reduced memory footprint
  - ▶ Faster algorithm since we avoid extra linear algebra overhead

## Lattice Ensembles

We use two RBC/UKQCD ensembles with physical mass Möbius DWF (MDWF) quarks:

- 1 **24ID**: Coarse simulation with periodic boundary conditions (RHMC or EOFA strange quark) [see Bob Mawhinney's talk, Tue. @ 17:30]
- 2 **32ID-G**: Production ensemble for  $\Delta I = 1/2 K \rightarrow \pi\pi$  calculation with  $G$ -parity boundary conditions (RHMC or EOFA light quarks) [Bai et al., arXiv:1505.07863]

Ensemble	$N_f$	Action	$\beta$	$L^3 \times T \times L_s$	Möbius Scale	GPBC	$am_l$	$am_h$
24ID	2+1	MDWF + ID	1.633	$24^3 \times 64 \times 24$	4.0	—	0.00789	0.085
32ID-G	2+1	MDWF + ID	1.75	$32^3 \times 64 \times 12$	2.67	$x,y,z$	0.0001	0.045

**Table:** Summary of ensembles and input parameters. Both ensembles use the Iwasaki + dislocation suppressing determinant ratio (DSDR) gauge action (ID). Here  $\beta$  is the gauge coupling, and “GPBC” denotes spatial directions with  $G$ -parity boundary conditions.

Ensemble	$a^{-1}$ (GeV)	$L$ (fm)	$m_\pi$ (MeV)	$m_K$ (MeV)
24ID	0.981(39)	4.82(19)	137.1(1)	494.6(1)
32ID-G	1.378(7)	4.57(2)	143.1(2.0)	490.6(2.4)

**Table:** Measured properties of both ensembles.

## EOFA Action [TWQCD, arXiv:1403.1683]

$$S_{\text{EOFA}} = \phi^\dagger \left[ \underbrace{\mathbb{1} - kP_- \Omega_-^\dagger [H(m_1)]^{-1} \Omega_- P_- + kP_+ \Omega_+^\dagger [H(m_2) - \Delta_+ P_+]^{-1} \Omega_+ P_+}_{\equiv \mathcal{M}_{\text{EOFA}}} \right] \phi$$

- $D_{\text{DWF}}$  (standard DWF Dirac operator) and  $D_{\text{EOFA}}$  (TWQCD's EOFA Dirac operator) are related by

$$D_{\text{DWF}} = D_{\text{EOFA}} \cdot \tilde{D}$$

where  $\tilde{D} = \delta_{xx'} \delta_{\alpha\alpha'} \delta_{aa'} (\tilde{D})_{ss'}$  is block diagonal in 4D spacetime, spin, and color indices, and has no gauge field dependence

- $H \equiv \gamma_5 R_5 D_{\text{EOFA}}$  is Hermitian EOFA Dirac operator
- Introducing  $D_{\text{EOFA}}$  trades explicit  $\gamma_5$ -Hermiticity for any choice of parameters for dense  $L_s \times L_s$  block structure in fifth dimension:  $D_{\text{DWF}}$  has tridiagonal  $ss'$  stencil but  $(\gamma_5 R_5 D_{\text{DWF}})^\dagger \neq \gamma_5 R_5 D_{\text{DWF}}$  in general
- Evaluating Hamiltonian or pseudofermion force requires two (ordinary) CG inversions
- Heatbath requires evaluating  $\phi = \mathcal{M}_{\text{EOFA}}^{-1/2} \eta$  (still need rational approximation!)

## Generic EOFA Linear System

$$(\gamma_5 R_5 D_{\text{EOFA}} + \beta \Delta_{\pm} P_{\pm}) \psi = \phi$$

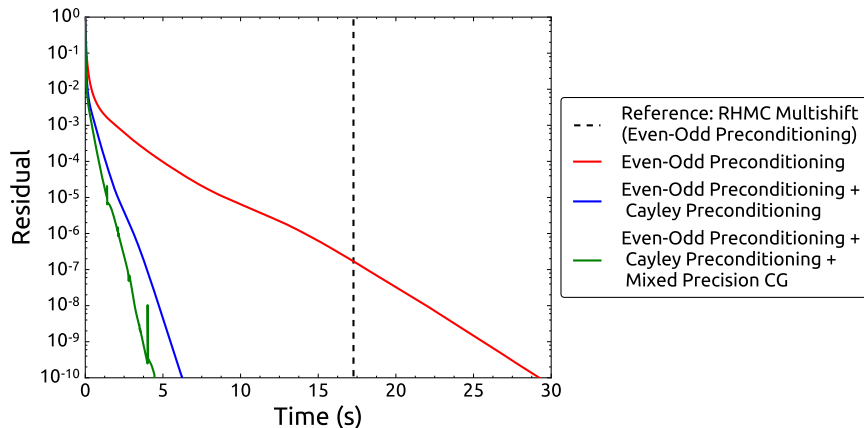
- Möbius  $D_{\text{EOFA}}$  is dense in  $ss'$  and thus expensive compared to (tridiagonal)  $D_{\text{DWF}}$
- Can use  $\tilde{D}^{-1}$  as a preconditioner to write EOFA system in terms of  $D_{\text{DWF}}$  for  $\beta = 0$ :

$$\boxed{\gamma_5 R_5 D_{\text{EOFA}} \psi = \phi} \iff \underbrace{D_{\text{EOFA}} \cdot \tilde{D}}_{=D_{\text{DWF}}} \cdot \underbrace{\tilde{D}^{-1} \psi}_{\equiv \psi'} = \gamma_5 R_5 \phi \iff \boxed{D_{\text{DWF}} \psi' = \phi'}$$

- $\beta \neq 0$  system can be treated as slight generalization of  $D_{\text{DWF}}$  to a four-point stencil in  $ss'$  by noticing that  $\Delta_{\pm} \tilde{D} = \vec{u} \otimes \vec{v}$  is rank-one [Jung et al., arXiv:1706.05843]
- **“Cayley preconditioning”**: solve substantially cheaper preconditioned system for  $\psi'$  and recover  $\psi$  for one additional matrix multiplication by  $\tilde{D}$
- Allows simple EOFA implementation reusing existing high-performance  $D_{\text{DWF}}$  code

## Cayley Preconditioning: 24ID Benchmark

**Benchmark:** Compare analogous solves required to compute the RHMC or EOFA Hamiltonian for the 24ID strange quark on a 256-node BG/Q partition



**Figure:** Wall clock time required to invert  $\gamma_5 R_5 D_{\text{EOFA}} \psi = \phi$  at the physical strange quark mass on the 24ID ensemble. The dashed vertical line corresponds to the total time required to invert  $(D_{\text{DWF}}^\dagger D_{\text{DWF}})^{1/2} \psi = \phi$  at the same mass using multishift CG.

## 24ID Evolution Benchmark

**Benchmark:** compare one molecular dynamics trajectory of the 24ID ensemble using either RHMC or EOFA for the strange quark

- ▶ Both simulations use identical force gradient (FG) integration schemes
- ▶ Can use FG forecasting [Yin and Mawhinney, arXiv:1111.5059] for EOFA evolution (ordinary CG) but not for RHMC evolution (multishift CG)
- ▶ Also compare EOFA with and without Cayley preconditioning

Step	RHMC		EOFA (Dense)		EOFA (Preconditioned)	
	Time (s)	%	Time (s)	%	Time (s)	%
Heatbath	42.6	2.7	340.6	15.1	68.9	15.5
Force gradient integration (total)	1485.6	94.8	1840.6	81.8	355.9	80.1
Final Hamiltonian evaluation	39.4	2.5	68.8	3.1	19.8	4.4
Total	1567.6	—	2250.0	—	444.6	—
<b>(Total RHMC) / Total</b>	<b>1.0</b>	—	<b>0.7</b>	—	<b>3.5</b>	—

**Table:** Strange quark timings for a single molecular dynamics (MD) trajectory of the 24ID ensemble on a 256-node BG/Q partition.

**3.5× speed-up in strange quark determinant  $\iff$  ~ 20% speed-up in job time**



## 32ID-G Evolution Tuning: Strategy

**First large-scale RBC/UKQCD EOFA calculation:** switch to EOFA for light ( $G$ -parity) quarks in ongoing  $\Delta I = 1/2 K \rightarrow \pi\pi$  calculation and re-tune integrator

- Previous RHMC run used four-level Omelyan integrator, with quark determinants split across two time steps, and a single Hasenbusch preconditioning mass for the light quark determinant [Bai et al., arXiv:1505.07863]

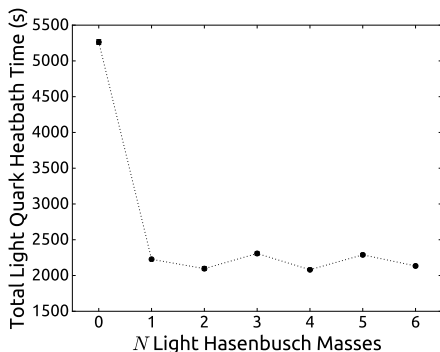
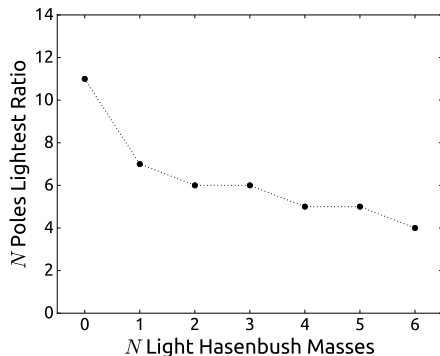
### Tuning Strategy

- 1 Switch to EOFA for the light quarks and three-level force gradient integrator with all quark determinants on same time step
- 2 Re-tune Hasenbusch mass preconditioning for (EOFA) light quark determinant
- 3 Tune heatbath rational approximations and CG tolerances
- 4 Re-tune step size and molecular dynamics CG tolerances

## 32ID-G Evolution Tuning: Hasenbusch Mass Preconditioning

In contrast to RHMC, Hasenbusch mass preconditioning works extremely well for EOFA!

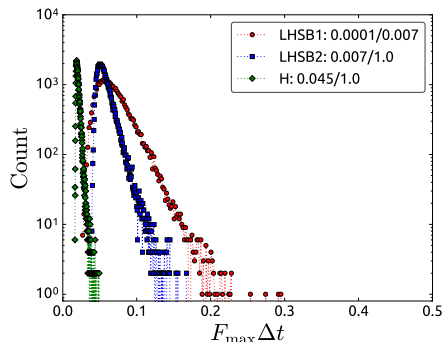
- Smaller pseudofermion forces  $\implies$  larger integration step size
- Also observe that spectral range of  $\mathcal{M}_{\text{EOFA}}$  contracts as  $m_2 \rightarrow m_1$ , allowing cheaper rational approximations to be used in the preconditioned heatbath



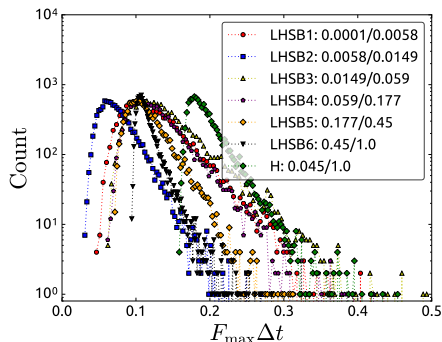
**Figure:** Left: minimum number of poles in the rational approximation to  $\mathcal{M}_{\text{EOFA}}^{-1/2}$  needed to achieve a relative error  $\epsilon < 10^{-10}$  for the lightest Hasenbusch ratio. Right: total light quark heatbath time including all Hasenbusch ratios.

## 32ID-G Evolution Tuning: Force Distributions

**Strategy:** Empirically observe that tails of force distributions control acceptance, so add intermediate Hasenbusch masses until light and heavy quark force distributions have comparable tails.



(a) RHC Ensemble



(b) EOFA Ensemble

**Figure:** Histograms of the maximum force,  $F_{\max} \equiv \max_{x,\mu} \|F_{\mu}(x)\|$ , measured on 850 trajectories of the 32ID-G RHC and EOFA ensembles. "LHSB" denotes the various mass ratios entering into our mass preconditioning scheme for the light quark determinant, and "H" denotes the strange quark determinant.

## 32ID-G Evolution Benchmark

Light Quark Action	Integrator	Light Hasenbusch Masses	$\Delta t$	$r_{\text{MD}}$	$N_{\text{traj}}$	Acceptance	Efficiency
RHMC	Omelyan	0.007	0.0625	$10^{-8}$	850	88%	—
EOFA	FG	0.0058, 0.0149, 0.059, 0.177, 0.45	0.1667	$10^{-7}$	850	93%	4.2

**Table:** Fully tuned RHMC and EOFA schemes.  $\Delta t$  is the outermost time step,  $r_{\text{MD}}$  is the MD CG tolerance, and **efficiency is the speed-up in the total job time relative to the RHMC scheme.**

### RHMC

- Omelyan integrator ( $\lambda = 0.22$ )
- One light quark Hasenbusch mass
- Multishift CG with single precision  $\mathcal{D}$  but accumulating solution and search vectors in double precision, coupled with reliable update to correct residual
- Even-odd preconditioning

### EOFA

- Force gradient integrator
- Five light quark Hasenbusch masses
- Mixed precision defect correction CG
- Even-odd preconditioning
- Cayley preconditioning
- Force gradient forecasting
- Heatbath forecasting
- Heatbath CG tolerance tuning

**Figure:** Comparison of RHMC and EOFA integration schemes.

## Conclusions

- We have introduced a new preconditioner for EOFA which relates inversions of (dense)  $D_{\text{EOFA}}$  to significantly cheaper inversions of (tridiagonal)  $D_{\text{DWF}}$
- We have also explored the performance of standard lattice preconditioners — even-odd, Hasenbusch — and additional tuning and forecasting techniques
  - ▶ EOFA appears much more amenable to Hasenbusch preconditioning than RHMC
  - ▶ No multishift CG in EOFA calculations means a larger set of tricks can be used (implicitly restarted solvers, forecasted CG initial guesses, ...)
- We find that physical quark mass Möbius DWF simulations can be substantially cheaper with EOFA compared to RHMC
  - ▶ 24ID ensemble:  $\sim 20\%$  speed-up in total job time per trajectory
  - ▶ 32ID-G ensemble:  $4.2\times$  speed-up in total job time per trajectory
- Ongoing  $\Delta I = 1/2 K \rightarrow \pi\pi$  ensemble generation is now being performed with EOFA evolution code [see Chris Kelly's talk, Wed. @ 13:10, for the physics!]
- Where else might EOFA be useful?
  - ▶  $N_f = 2 + 1 + 1$  simulations with dynamical (EOFA) strange and charm quarks
  - ▶ Simulations with physical, isospin-broken light quarks (hadron spectra,  $K_{\ell 3}$ , ...)
  - ▶ Simulations with  $SU(3)$ -symmetric light quarks ( $\chi$ PT studies, thermodynamics, ...)

Thank you!



# Backup Slides

## Refinement II: Heatbath Tuning

- EOFA heatbath: draw Gaussian field  $\eta$  and seed  $\phi = \mathcal{M}_{\text{EOFA}}^{-1/2} \eta$
- Can show

$$\mathcal{M}_{\text{EOFA}}^{-1/2} \simeq \alpha_0 + \sum_{k=1}^{N_p} \alpha_l \gamma_l \left\{ \mathbb{1} + k\gamma_l P_- \Omega_-^\dagger [H(m_1) - \gamma_l \Delta_-(m_1, m_2) P_-]^{-1} \Omega_- P_- \right. \\ \left. - k\gamma_l P_+ \Omega_+^\dagger [H(m_2) - \beta_l \gamma_l \Delta_+(m_1, m_2) P_+]^{-1} \Omega_+ P_+ \right\}$$

for a rational approximation  $x^{-1/2} \simeq \alpha_0 + \sum_l \alpha_l / (\beta_l + x)$ , with  $\gamma_l \equiv (1 + \beta_l)^{-1}$

- Requires  $2N_p$  independent CG solves, since  $\Delta_\pm$  operators are singular  $\rightarrow$  expensive!
- Cost can be partially ameliorated by forecasting initial guesses across poles [TWQCD, arXiv:1403.1683; Murphy, arXiv:1611.00298]
- Further reduction possible by observing that coefficients  $\{\alpha_0, \alpha_l \gamma_l, k\alpha_l \gamma_l^2\}$  generally exhibit large separation of scales, suggesting loose CG tolerances can be used for inversions multiplied by small coefficients



## Refinement II: Heatbath Tuning

We propose tuning the heatbath CG tolerances by considering the quantity

$$\varepsilon \equiv \frac{|\eta^\dagger \eta - \phi^\dagger \mathcal{M}_{\text{EOFA}} \phi|}{\eta^\dagger \eta} \quad \left( \phi = \mathcal{M}_{\text{EOFA}}^{-1/2} \eta \right)$$

and using the following simple algorithm:

- 1 Choose a desired overall tolerance for the heatbath,  $\varepsilon_{\text{tol}}$
- 2 Choose one of the inversions required to compute  $\phi$ , and relax the stopping tolerance until the overall error in the heatbath ( $\varepsilon$ ) reaches  $\varepsilon_{\text{tol}}$
- 3 Iterate over each inversion until all stopping conditions have been tuned

Ensemble		$\varepsilon$	Total Heatbath Time	Speed-Up
24ID	Untuned	$1.52 \times 10^{-11}$	129.5 s	<b>1.9</b>
	Tuned	$7.79 \times 10^{-11}$	68.9 s	
32ID-G	Untuned	$1.11 \times 10^{-11}$	2289.0 s	<b>2.7</b>
	Tuned	$7.69 \times 10^{-11}$	840.1 s	

Table: Relative error ( $\varepsilon$ ) and total running time for the EOFA heatbath before and after tuning.

## 24ID and 32ID-G Integrator Layouts

Ensemble	Level	Action	Update
24ID (RHMC)	1	$\text{RatQuo}_{1/2}(0.085, 1.0)$	1:1
	2	$\text{Quo}(0.00107, 0.00789) + \text{Quo}(0.00789, 0.0291) + \text{Quo}(0.0291, 0.095) +$ $\text{Quo}(0.095, 0.3) + \text{Quo}(0.3, 0.548) + \text{Quo}(0.548, 1.0)$	1:1
	3	Gauge + DSDR	1:1
24ID (EOFA)	1	$\text{EOFA}(0.085, 1.0)$	1:1
	2	$\text{Quo}(0.00107, 0.00789) + \text{Quo}(0.00789, 0.0291) + \text{Quo}(0.0291, 0.095) +$ $\text{Quo}(0.095, 0.3) + \text{Quo}(0.3, 0.548) + \text{Quo}(0.548, 1.0)$	1:1
	3	Gauge + DSDR	1:1
32ID-G (RHMC)	1	$\text{RatQuo}_{1/2}(0.0001, 0.007)$	1:1
	2	$\text{RatQuo}_{1/2}(0.007, 1.0) + \text{RatQuo}_{1/4}(0.045, 1.0)$	1:2
	3	DSDR	1:2
	4	Gauge	1:1
32ID-G (EOFA)	1	$\text{EOFA}(0.0001, 0.0058) + \text{EOFA}(0.0058, 0.0149) + \text{EOFA}(0.0149, 0.059) +$ $\text{EOFA}(0.059, 0.177) + \text{EOFA}(0.177, 0.45) +$ $\text{EOFA}(0.45, 1.0) + \text{RatQuo}_{1/4}(0.045, 1.0)$	5:1
	2	DSDR	1:2
	3	Gauge	1:1

**Table:** Integrator layouts for the 24ID and 32ID-G ensembles. The notation A:B for the update scheme denotes the number of steps of the next innermost integrator level (A) per step of the current level (B).